

Lecture Summary For Friday May 1st

We looked at the shallow water system again with Beta turned off, and we introduced a wall into the system. We also assumed that $u' = 0$ everywhere in the system and for $t =$ any

number. As a result, our shallow water equations are reduced to: $f_0 \dot{v} = -g \frac{d\hat{h}}{dx}$, $\frac{d\dot{v}}{dt} = -g \frac{d\hat{h}}{dy}$,

and $\frac{d\hat{h}}{dt} + H \frac{d\dot{v}}{dy} = 0$. Since $u' = 0$, we see that the wave has to propagate in the y -direction. If we

take the partial with respect to “ t ” of the third shallow water equation we are left with

$\frac{d^2 \hat{h}}{dt^2} + H \frac{d}{dy} \left(\frac{d\dot{v}}{dt} \right) = 0$. Then we took the partial with respect to “ y ” of the second equation in

order to make a substitution. We are left with the differential equation: $\frac{d^2 \hat{h}}{dt^2} = gH \frac{d^2 \hat{h}}{dy^2}$. From

this equation we are able to get the expression $gH = C_{sw}^2$ where C is the shallow water

propagation of the wave. Also from this differential equation we get the solutions

$\hat{h} = f(y + C_{sw} t)$ for southward propagation and $\hat{h} = f(y - C_{sw} t)$.

Then we plugged $\hat{h} = \tilde{h}(y) \exp(i(\omega t - ky))$ into our differential equation which we then

simplified to $\frac{\omega^2}{k^2} = gH$. From this relationship, we were able to make the equality

$$C_y = \pm \sqrt{gH}$$